

## NMR and NQR study of pressure-induced superconductivity and the origin of critical-temperature enhancement in the spin-ladder cuprate $\text{Sr}_2\text{Ca}_{12}\text{Cu}_{24}\text{O}_{41}$

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(Received 6 August 2009; published 8 September 2009)

Pressure-induced superconductivity was studied for a spin-ladder cuprate  $\text{Sr}_2\text{Ca}_{12}\text{Cu}_{24}\text{O}_{41}$  using nuclear magnetic resonance under pressures up to the optimal pressure 3.8 GPa. Pressure application leads to a transitional change from a spin-gapped state to a Fermi-liquid state at temperatures higher than  $T_c$ . The relaxation rate  $1/T_1$  shows activated-type behavior at an onset pressure, whereas Korringa-like behavior becomes predominant at the optimal pressure, suggesting that an increase in the density of states at the Fermi energy leads to enhancement of  $T_c$ . Nuclear quadrupole resonance spectra suggest that pressure application causes transfer of holes from the chain to the ladder sites. The transfer of holes increases DOS below the optimal pressure. A dome-shaped  $T_c$  versus pressure curve arises from naive balance between the transfer of holes and broadening of the band width.

DOI: [10.1103/PhysRevB.80.100503](https://doi.org/10.1103/PhysRevB.80.100503)

PACS number(s): 74.25.Ha, 74.62.Fj, 74.70.-b, 74.72.Jt

A spin-ladder system,  $\text{Sr}_{14-x}\text{Ca}_x\text{Cu}_{24}\text{O}_{41}$  ( $x=11.5-13.5$ ), which is known as the only cuprate superconductor with nonsquare lattice Cu-O layers, offers a good model for study of quasi-one-dimensional superconductivity. To date, various theoretical studies have been performed for isolated two-leg ladders or the Trellis lattice.<sup>1-11</sup> Although the system is attractive for comparison with high- $T_c$  cuprates, little experimental effort has been devoted to investigation of superconductivity because of experimental difficulties under high pressure; the onset and optimal pressures are around 3 and 4 GPa, respectively.<sup>12,13</sup> The experiments at high pressures, above 3 GPa, have been performed only for resistivity, ac susceptibility, x-ray diffraction and NMR measurements.<sup>12-21</sup> Appearance of a dome-shaped  $T_c$  curve on the  $P$ - $T$  phase diagram and a decrease in the spin gap caused by pressure application are reminiscent of hole doping in high- $T_c$  cuprates.<sup>13,18,19</sup> The pressure-induced superconducting state survives even at high fields beyond the Pauli paramagnetic limit, when the field  $\mathbf{H}$  is applied parallel to the leg direction ( $\mathbf{H}\parallel\mathbf{c}$ ), for which orbital suppression is depressed.<sup>15</sup> The NMR measurement at 3.5 GPa shows that the quasiparticle excitation accompanies a full gap; however, the Knight shift shows no appreciable change below  $T_c$ .<sup>20</sup>

This system also shows features common to quasi-one-dimensional systems such as organic superconductors. An example is emergence of a charge-density wave (CDW) in a lightly doped regime ( $x\leq 0.8$ ) at ambient pressure.<sup>22,23</sup> The stable superconducting state beyond the Pauli paramagnetic limit and no change in the Knight shift below  $T_c$  are also common to  $(\text{TMTSF})_2\text{X}$  ( $\text{X}=\text{ClO}_4$  or  $\text{PF}_6$ ). These features have often been discussed in relation to the possibility of  $p$ -wave superconductivity or singlet-Fulde-Ferrell-Larkin-Ovchinnikov (FFLO) phase transition.<sup>24-26</sup> Recently, in the cuprate spin-ladder system, the possibility of FFLO state has been investigated theoretically at high field.<sup>8-10</sup> NMR study at optimal pressure is required to investigate the reasons that

lead to  $T_c$  enhancement after pressure application. Questions that need to be addressed are: (1) what is the origin of the dome-shaped  $T_c$  curve and (2) what is the spin susceptibility ( $\chi_{spin}$ ) like around  $T_c$ . We performed  $^{63}\text{Cu}$ -NMR measurements up to the optimal pressure ( $\sim 3.8$  GPa) to study these problems.

NMR and NQR measurements were performed using a single crystal of  $\text{Sr}_2\text{Ca}_{12}\text{Cu}_{24}\text{O}_{41}$  having a volume of  $4\times 2\times 1\text{ mm}^3$ . The measurements up to 3.5 GPa were carried out using a clamped-type pressure cell. A special variable-temperature insert (VTI) equipped with a 45 ton press was used to achieve a high pressure of 3.8 GPa. Details of the apparatus are described in Fujiwara *et al.*<sup>27</sup> We also determined  $T_c$  from resistivity measurements using a cubic anvil cell. A top-loading insert with a 250 ton press was used to control a steady load at low temperatures. Details of the apparatus and calibration method are described in Mōri *et al.*<sup>28</sup>

Figure 1 shows  $1/T_1$  of  $^{63}\text{Cu}$  nuclei measured at 6.2 T for  $\mathbf{H}\parallel\mathbf{c}$  up to 3.8 GPa. The rate  $1/T_1$  shows activated-type  $T$  dependence, which originates from the spin-gap excitation, followed by  $T$ -linear dependence;

$$T_1^{-1} = Ae^{-\Delta/T} + BT. \quad (1)$$

The superconducting state appears at low temperatures accompanied by a hump, implying existence of a full gap for the quasiparticle excitation. The coefficient of the  $T$ -linear dependence [B in Eq. (1)] is enhanced with increasing pressure up to 3.8 GPa, similar to  $T_c$ . The enhancement of  $T_c$  is confirmed by a shift of the coherence peak to a higher temperature. The  $1/T_1T$  vs  $T$  plot normalized by those quantities at  $T_c$  fits a single curve. The  $T$ -linear dependence comes from the Korringa relation, and thus, the coefficient  $B$  is proportional to the square of the density of states (DOS) at the Fermi energy,  $D(E_F)$ . Therefore, the enhancement of  $T_c$  is attributable to an increase in  $D(E_F)$ . On the other hand, the

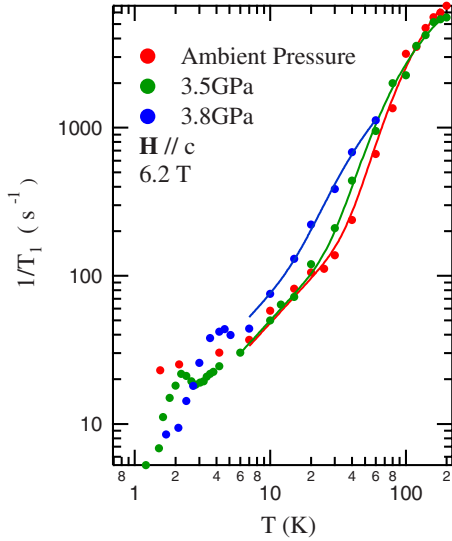


FIG. 1. (Color online) Nuclear magnetic relaxation rate  $1/T_1$  of  $^{63}\text{Cu}$  nuclei measured at 6.2 T for  $\mathbf{H}\parallel c$ , the leg direction. The data are fitted by Eq. (1).

spin gap  $\Delta$  and coefficient  $A$  in Eq. (1) decrease with increasing pressure. The values of  $A$  at ambient pressure, 3.5 and 3.8 GPa are in a ratio of about 6:4:1. The activated-type contribution fades out with increasing pressure, although both mechanisms coexist in the relevant pressure regime. The rate  $1/T_1$  under extremely high pressure is expected to become  $1/T_1 T = \text{constant}$ . The  $P$  dependence of both  $B$  and  $\Delta$  is shown in Fig. 2, together with that of  $T_c$  determined from the resistivity measurements. The data determined from the onset and zero resistivity are plotted in Fig. 2. We also plotted the data of  $T_c$  determined from the inductance of the NMR probe. Our results show that  $\Delta$  is almost the same up to the onset pressure, but decreases drastically above it. Corresponding to  $\Delta$ ,  $B$  shows a remarkable upturn toward the optimal pressure. The upturn of  $B$  ( $=T$  independent) means that antiferromagnetic spin fluctuation hardly affects the en-

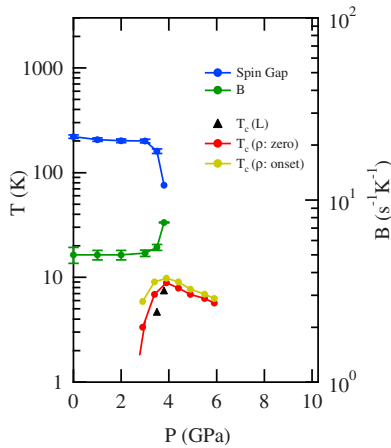


FIG. 2. (Color online) Pressure dependence of spin gap  $\Delta$  and the coefficient of  $T$ -linear term,  $B$  in Eq. (1). Closed triangles represent  $T_c$  determined from the inductance of the NMR probe.  $T_c(\rho: \text{onset})$  and  $T_c(\rho: \text{zero})$  represent values of  $T_c$  determined from the onset and zero resistivity, respectively.

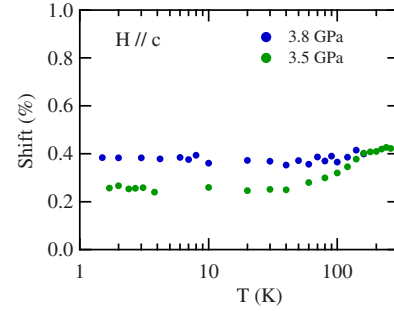


FIG. 3. (Color online) Knight shift of  $^{63}\text{Cu}$  nuclei for the ladder sites at low temperatures derived from high fields.

hancement of  $T_c$ ; when antiferromagnetic fluctuations are predominant such as high- $T_c$  cuprates,  $T$  dependence of  $1/T_1 T$  shows a upturn toward  $T_c$ .

An increase in  $D(E_F)$  is also observed from the Knight shift for  $\mathbf{H}\parallel c$ ; the spin-gap behavior observed clearly at 3.5 GPa almost vanishes at 3.8 GPa (see Fig. 3). The Knight shift is decomposed into the orbital and spin parts;  $K^c = K_{orb}^c + K_{spin}^c$ . The orbital part  $K_{orb}^c$  is almost unchanged for a pressure of 3.5–3.8 GPa as explained in the following paragraphs. Therefore, the spin part  $K_{spin}^c$  or  $\chi_{spin}$ , which is proportional to  $D(E_F)$  in a normal metal, is enhanced with increasing pressure. We estimated  $K_{orb}^c$  from the NQR measurements because  $K_{orb}^c$  is related to the NQR frequency  $\nu$  as

$$K_{orb}^c = \alpha\nu + \beta. \quad (2)$$

The relation is derived from Eqs. (3)–(5) in the following paragraphs.

Before presenting the results, we state the NQR results measured at 4.2 K. The NQR spectra at 3.5 GPa and ambient pressure are shown in Fig. 4. We measured the NQR signals up to 40 MHz with a resolution of 0.1 MHz/point. The signals originating from the ladder and chain sites were observed at 10–25 and 30–35 MHz, respectively. At ambient pressure,  $^{63}\text{Cu}$  and  $^{65}\text{Cu}$  signals are clearly assigned only in case of chain sites, whereas they overlap with each other in case of ladder sites. The overlap is decomposed into two Gaussian functions. The resonance position is almost the same as that measured in an earlier study.<sup>29</sup> At 3.5 GPa, the signals arising from the ladder sites move to high frequencies, whereas those from the chain sites move to low frequencies. The broadening of the line width is caused by

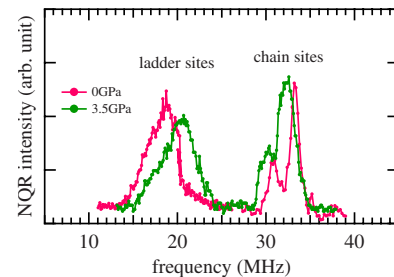


FIG. 4. (Color online) NQR spectra measured at 4.2 K for the ladder (15–25 MHz) and chain sites (30–35 MHz).

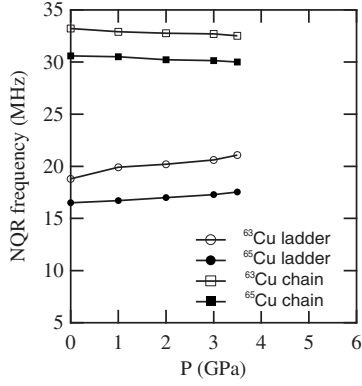


FIG. 5. Pressure dependence of the NQR frequency of  $^{63}\text{Cu}$  and  $^{65}\text{Cu}$  nuclei.

stress anisotropy, which appears when pressure mediation liquid freezes during the pressurizing process. The  $P$  dependence of the NQR frequencies is shown in Fig. 5. The  $^{63}\text{Cu}$  signals move by 2.27 and  $-0.69$  MHz for the ladder and chain sites, respectively, and those of  $^{65}\text{Cu}$  move by 1.05 and  $-0.60$  MHz for the ladder and chain sites, respectively. The  $\nu$  in Eq. (2) is determined by two factors; the electric-field gradient (EFG) arising from the hole number of the  $3d_{x^2-y^2}$  orbital,  $n_{3d}$  and that arising from the surrounding  $4p$  orbitals. In high- $T_c$  cuprates, the latter contribution is not so sensitive to doping level; therefore,  $\nu$  is determined mainly by  $n_{3d}$ ,

$$\nu = n_{3d}\nu_{3d} + \nu_{4p} \quad (3)$$

where  $\nu_{4p} \approx -65$  MHz and  $\nu_{3d} \approx -117$  MHz are estimated.<sup>30</sup> Equation (3) can be applied to both ladder and chain sites of  $\text{Sr}_{14-x}\text{Ca}_x\text{Cu}_{24}\text{O}_{41}$  as well as high- $T_c$  cuprates. The increase in the number of holes at the ladder Cu sites comes from the chain sites. We estimated  $n_{3d}$  from the resonance frequencies of  $^{63}\text{Cu}$  as 0.716 and 0.736 at ambient pressure and 3.5 GPa, respectively.

The hole number is related to the orbital susceptibility  $\chi_{orb}$ , namely,  $K_{orb}^c$ . The orbital susceptibility for  $\mathbf{H}\parallel c$  is expressed in the same way with high- $T_c$  cuprates<sup>31</sup> as

$$\frac{\chi_{orb}}{N} = n_{3d} \frac{2\mu_B^2}{E_{xy} - E_{x^2-y^2}} \quad (4)$$

where  $N$  is the Avogadro number and  $\mu_B$  is the Bohr magneton.  $E_{xy}$  and  $E_{x^2-y^2}$  represent energy levels of the  $3d_{xy}$  and

$3d_{x^2-y^2}$  orbitals, respectively. The energy difference between them is estimated to be 4.8 eV from the cluster model.<sup>32</sup>  $K_{orb}^c$  is linked with  $\chi_{orb}$  as

$$K_{orb}^c = A^{orb} \chi_{orb} / N \mu_B \quad (5)$$

where  $A^{orb} (\equiv 2\mu_B \langle 1/r^3 \rangle)$ ,  $\langle 1/r^3 \rangle$  being a radius average for a Cu ion, is a hyperfine field due to the orbital moment induced by the applied field.  $A^{orb}$  is estimated as 1028 (kOe/ $\mu_B$ ) using  $\langle 1/r^3 \rangle = 8.252$  a.u., a value for a free  $\text{Cu}^{2+}$  ion.<sup>33</sup> From Eqs. (4) and (5),  $K_{orb}^c = 0.248n_{3d}(\%)$  is obtained. A  $K_{orb}^c$  of 0.18% is estimated in a pressure of 3.5–3.8 GPa. The estimate shows that  $K_{spin}^c$  is enhanced at the optimal pressure.

The increase in  $T_c$  or  $D(E_F)$  is determined by two factors; the hole number and the band width. The hole number seems to increase beyond the optimal pressure, as is expected from Fig. 5, whereas the band width tends to broaden with increasing pressure. Therefore, the dome-shaped  $T_c$  curve arises from naive balance between the two factors: the transfer of holes to the ladder sites increases DOS below the optimal pressure, whereas the broadening of the band width would cause a decrease in DOS over the optimal pressure.

Anomalous behavior, namely, that  $K_{spin}^c$  is finite below  $T_c$ , is actually observed at rather high fields. If it is related to quasione dimensionality, the singlet-FFLO phase transition is a candidate that can explain the anomaly. In this scenario,  $K_{spin}^c$  or  $\chi_{spin}$  at a low field would decrease with decreasing temperature below  $T_c$ . Although NMR measurements at a low field is important, it is difficult to carry out because the overlapped NQR signals in Fig. 4 hardly split at a low field.

In summary, we have measured  $1/T_1$  at an optimal pressure of 3.8 GPa, and observed transitional features from the spin-gapped state to the Fermi-liquid state; the activated-type contribution fades out and approaches the Korringa relation with increasing pressure. This implies that antiferromagnetic spin fluctuation is not important unlike high- $T_c$  cuprates. The enhancement of  $T_c$  or an increase in DOS at the Fermi energy originates from an increase in the number of holes transferred from the chain sites.

The authors wish to thank S. Fujimoto and K. Kojima for their fruitful discussions and K. Tatsumi and A. Hisada for their experimental support. The present work was partially supported by Grant-in-Aid (Grant No. KAKENHI 17340107) for the Ministry of Education, Science and Culture, Japan.

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